



Relativistic Electron Loss Timescales in the Slot Region

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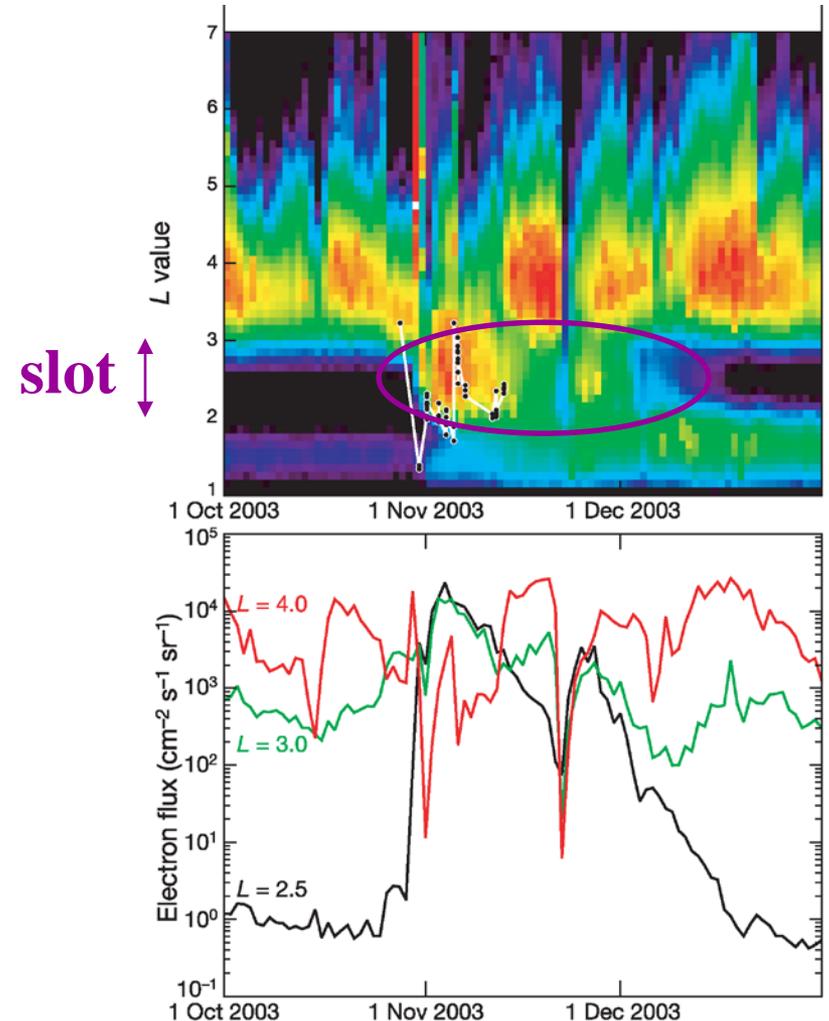
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Slot Region Dynamics

- Slot region
 - is not always empty.
 - can become filled during exceptionally large storms such as the Halloween Storms of 2003.
- subsequently reforms over the following weeks to months.

SAMPEX 2-6 MeV Electrons



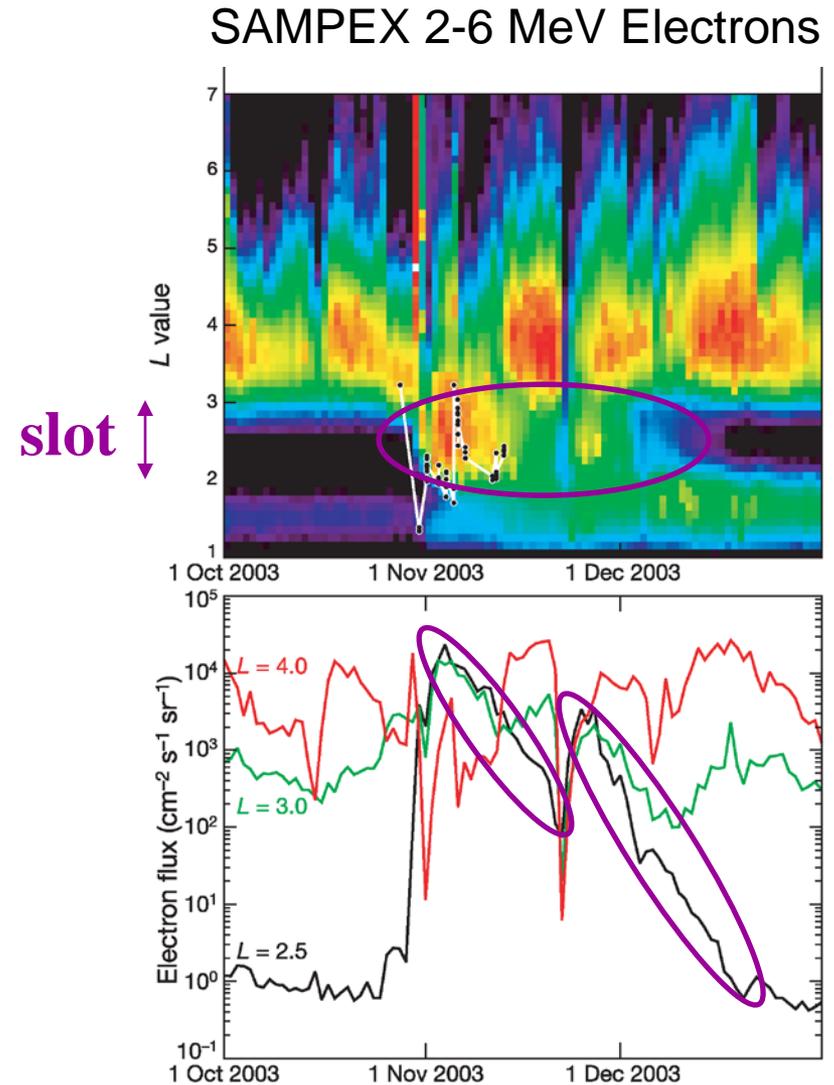
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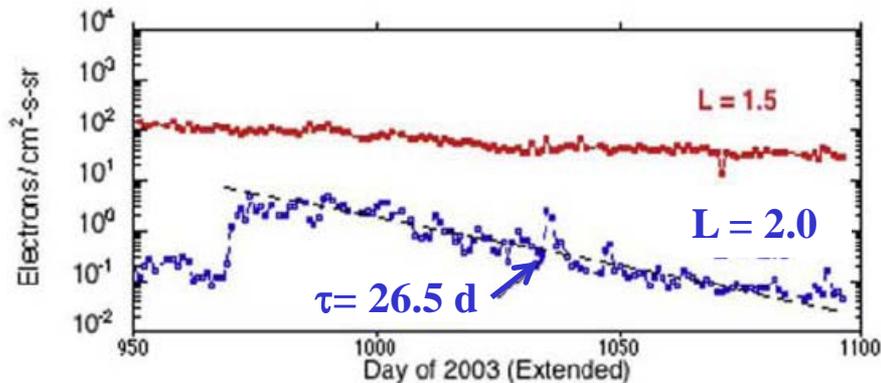
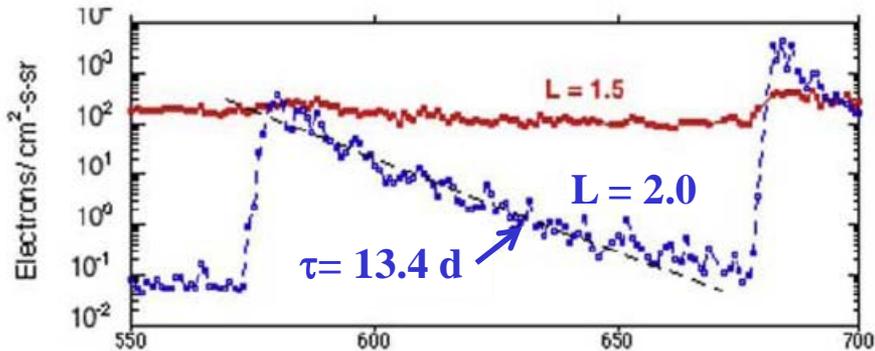
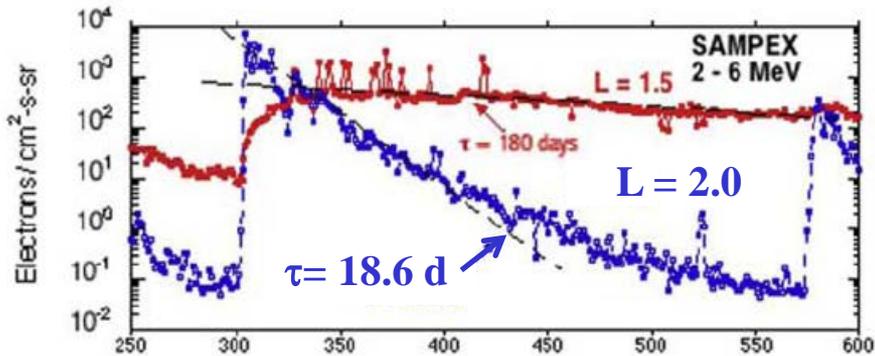
Baker *et al.*, Nature, 2004

Slot Region Loss Timescales

- Loss timescales for 2-6 MeV electrons in the centre of the slot at $L = 2.5$ are estimated to be of the order of 2.9 – 4.6 days.
- This is consistent with theoretical expectations based on losses due to plasmaspheric hiss [[Meredith *et al.*, 2007](#)].

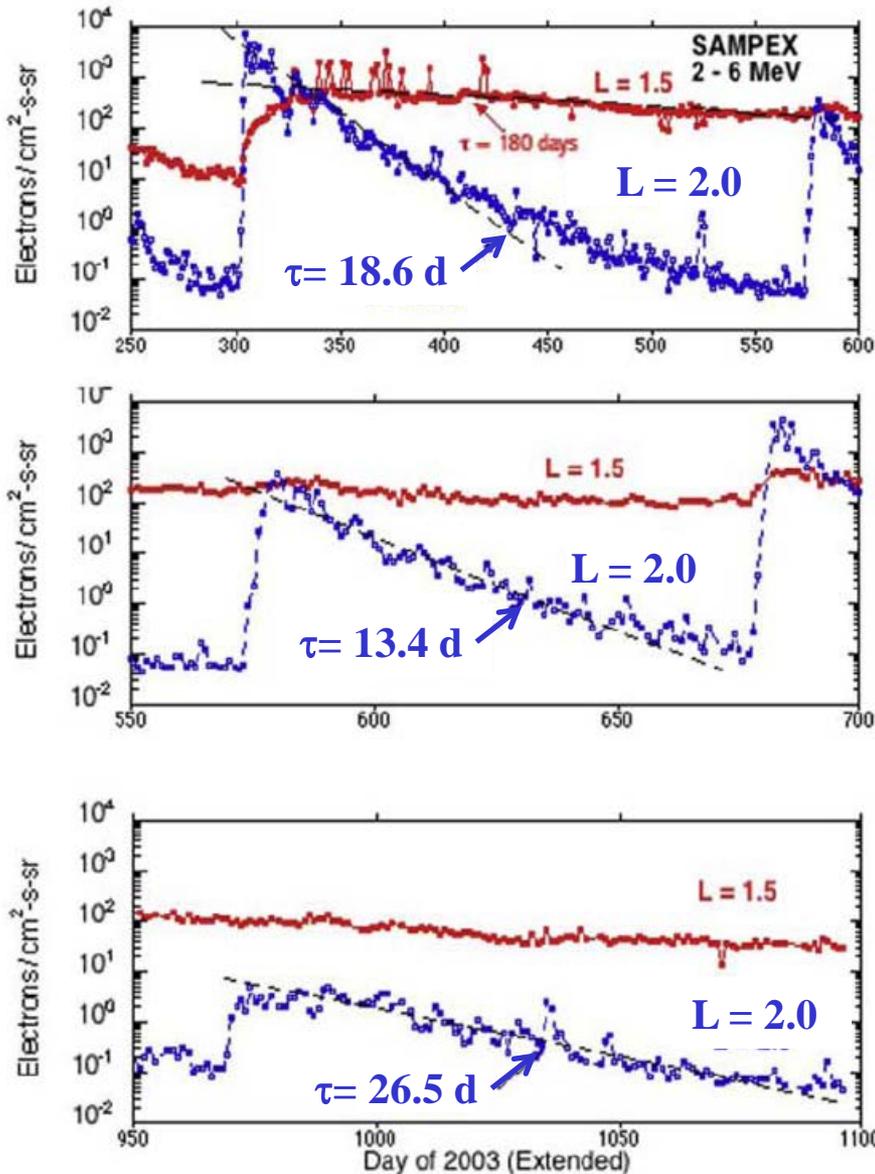


Slot Region Loss Timescales



- The experimental lifetime at L = 2.0 is ~20 days [[Baker et al., 2007](#)].
- This lifetime is much shorter than the theoretical estimates of a few hundred days as a result of losses due to plasmaspheric hiss alone. [[Meredith et al., 2007](#)].

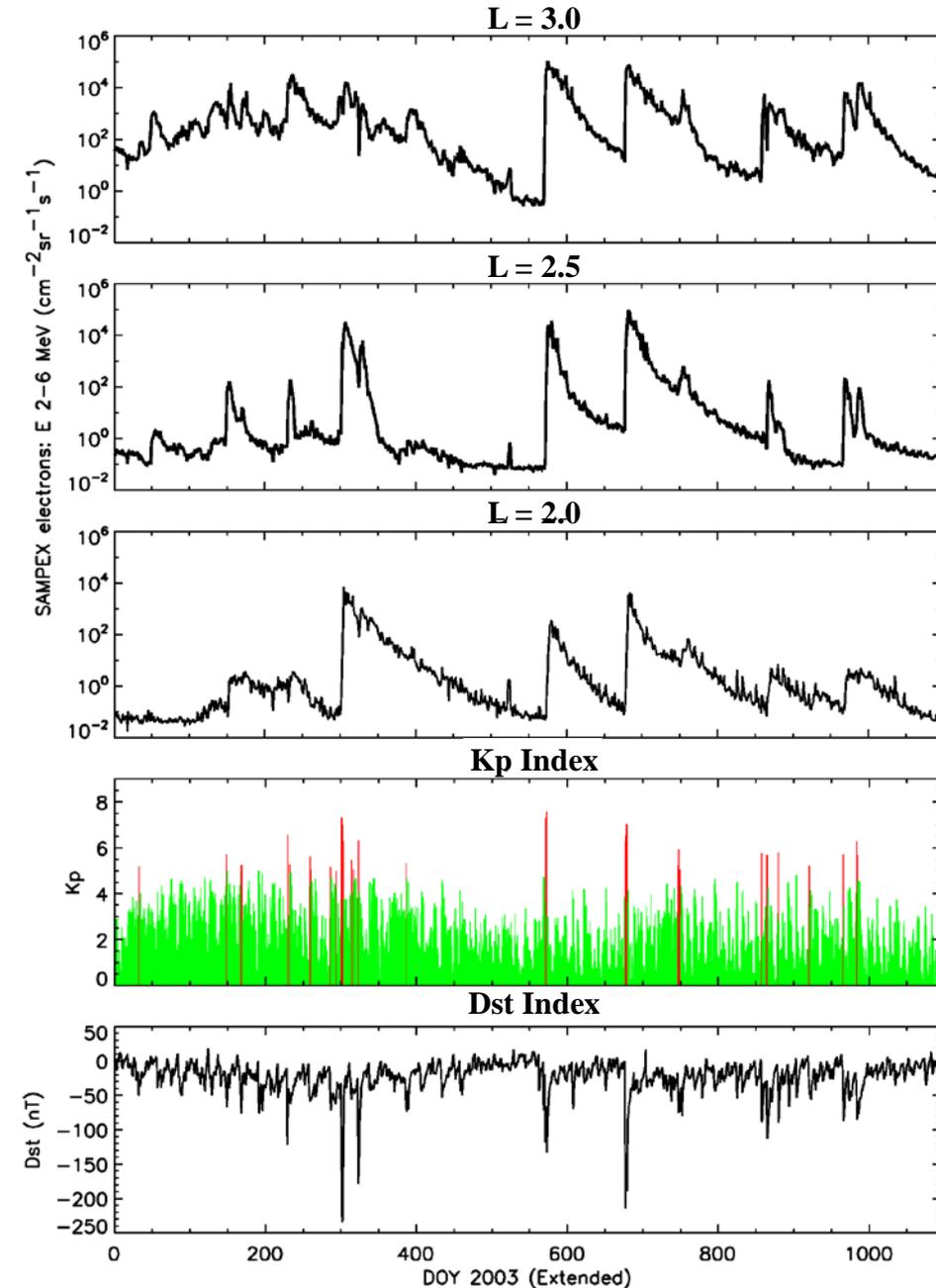
Slot Region Loss Timescales



- In this study we perform a comprehensive survey of relativistic electron loss timescales in the slot region.
- We compare SAMPEX observations with theoretical predictions derived from wave models based on CRRES observations.

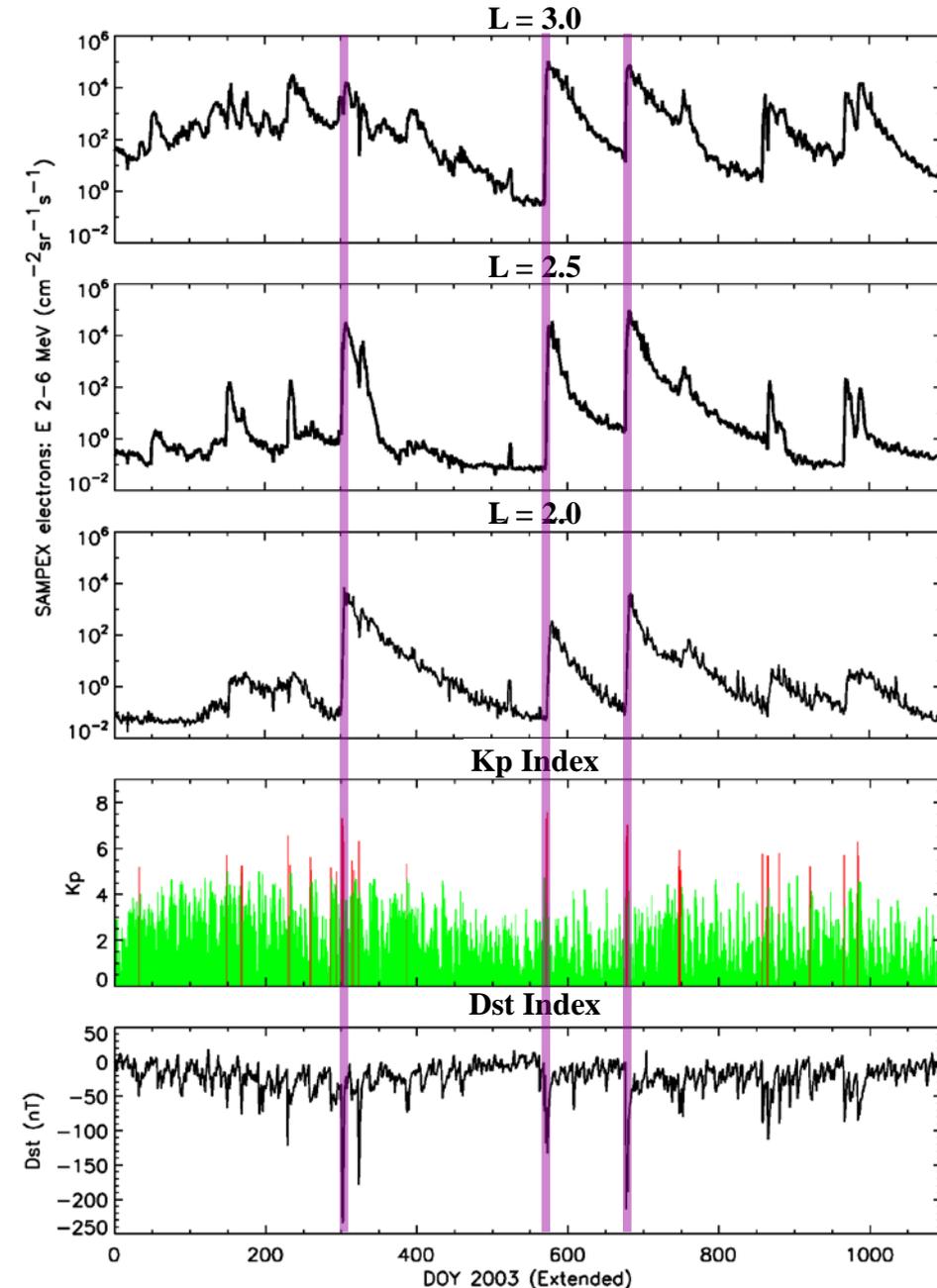
Slot Region Dynamics

- The flux at each location is characterised by a number of abrupt rises followed by gradual exponential decay.
- The abrupt increases
 - can be as large as 5 orders of magnitude.



Slot Region Dynamics

- The flux at each location is characterised by a number of abrupt rises followed by gradual exponential decay.
- The abrupt increases
 - can be as large as 5 orders of magnitude.
 - are associated with strong storms and enhanced magnetic activity.

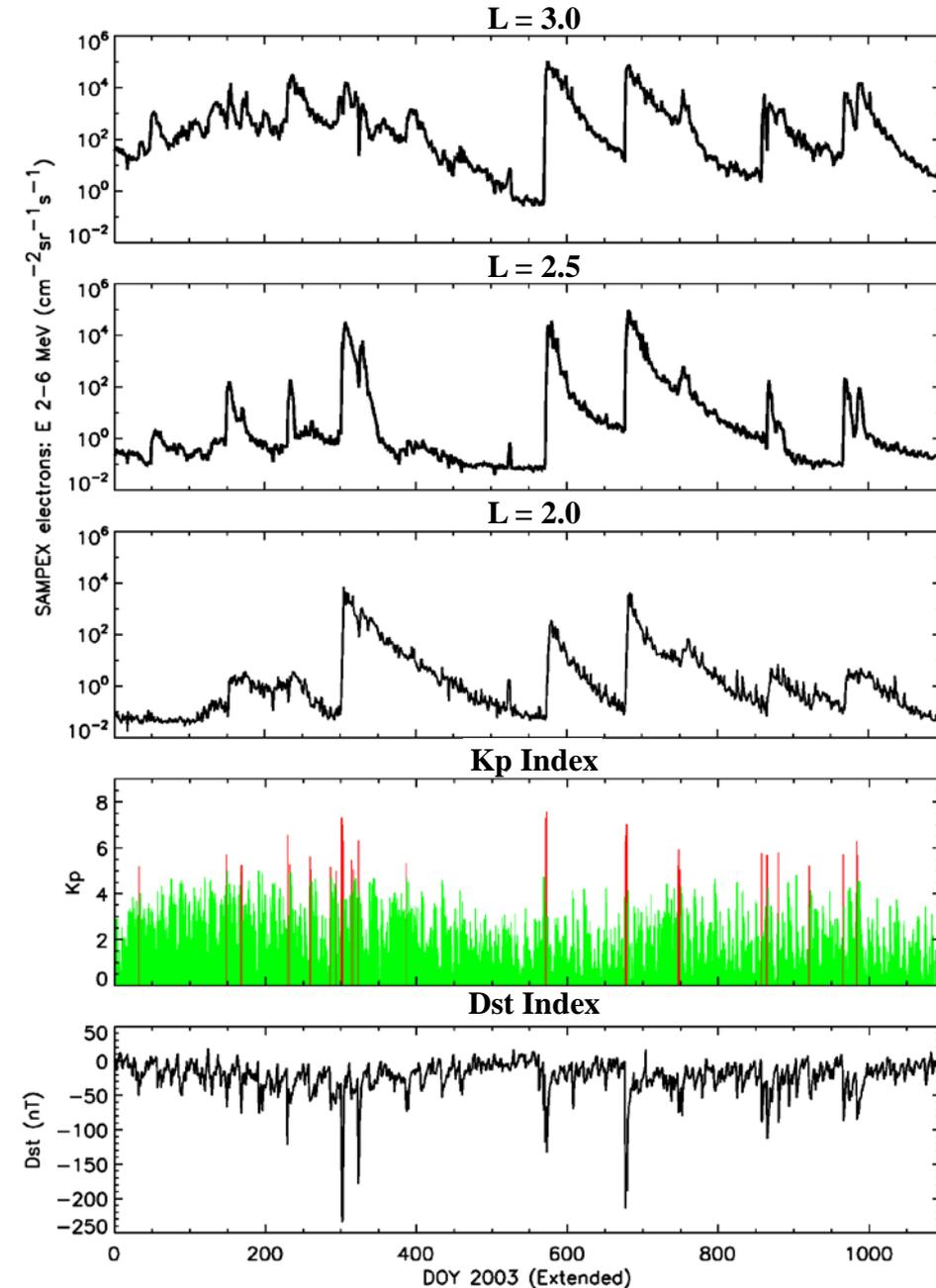


Loss Timescales

- To quantify the decay rates we fit an exponential function of the form:

$$J = J_0 \exp(-t/\tau)$$

- where τ is the loss timescale

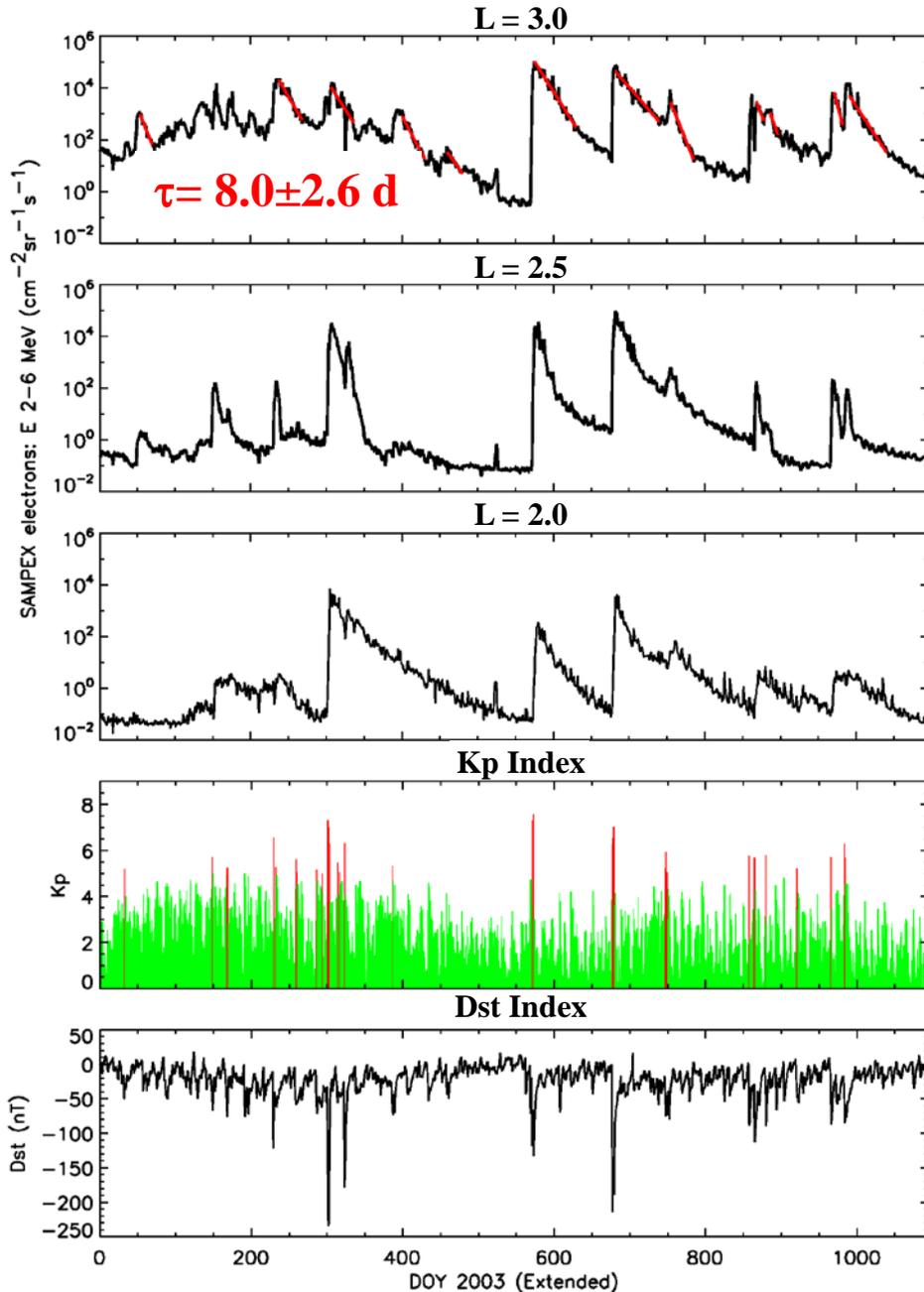


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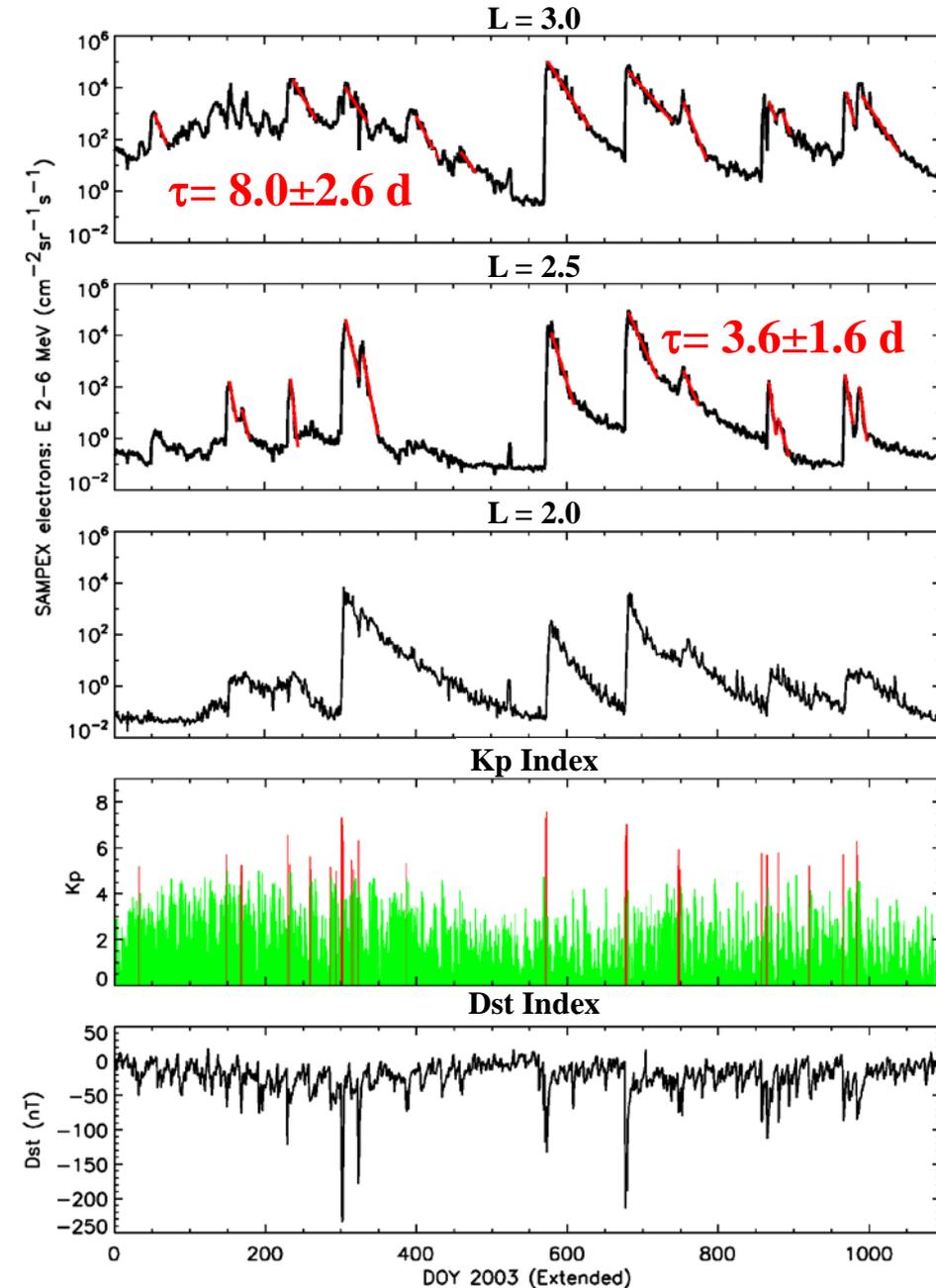
L	τ (days)
3.0	8.0 ± 2.6

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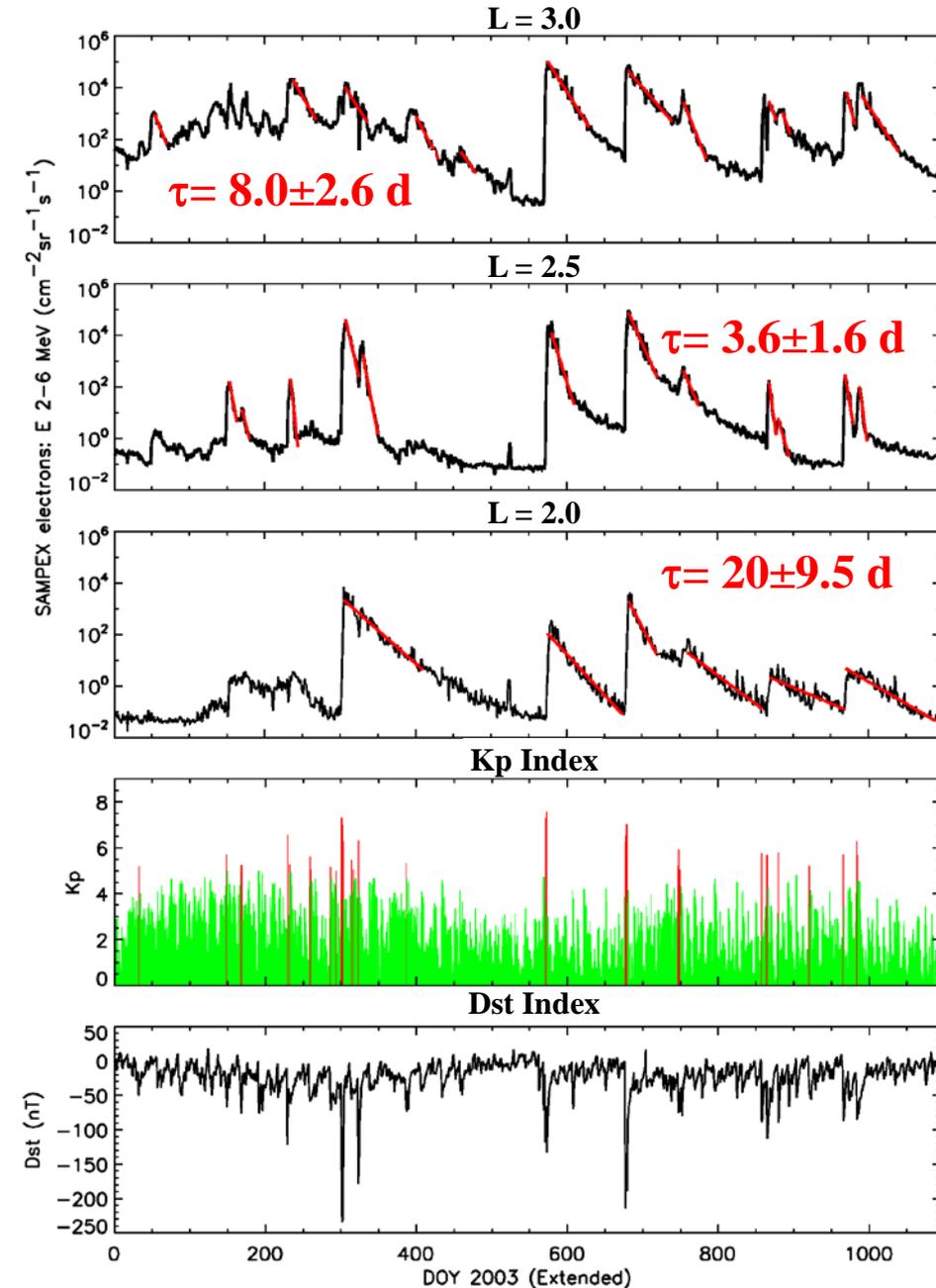
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3.0	8.0 ± 2.6
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Loss Timescales

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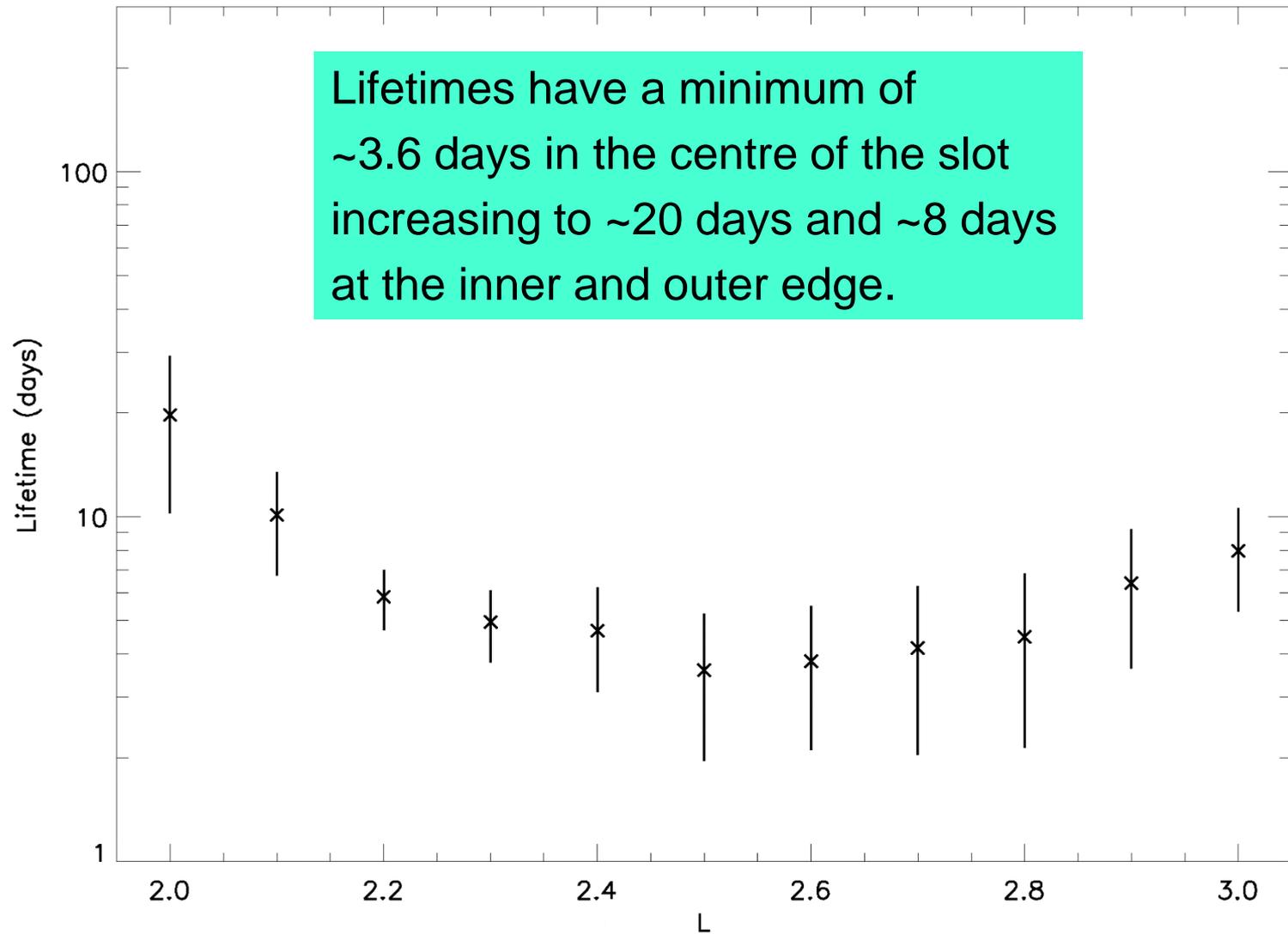
$$J = J_0 \exp(-t/\tau)$$

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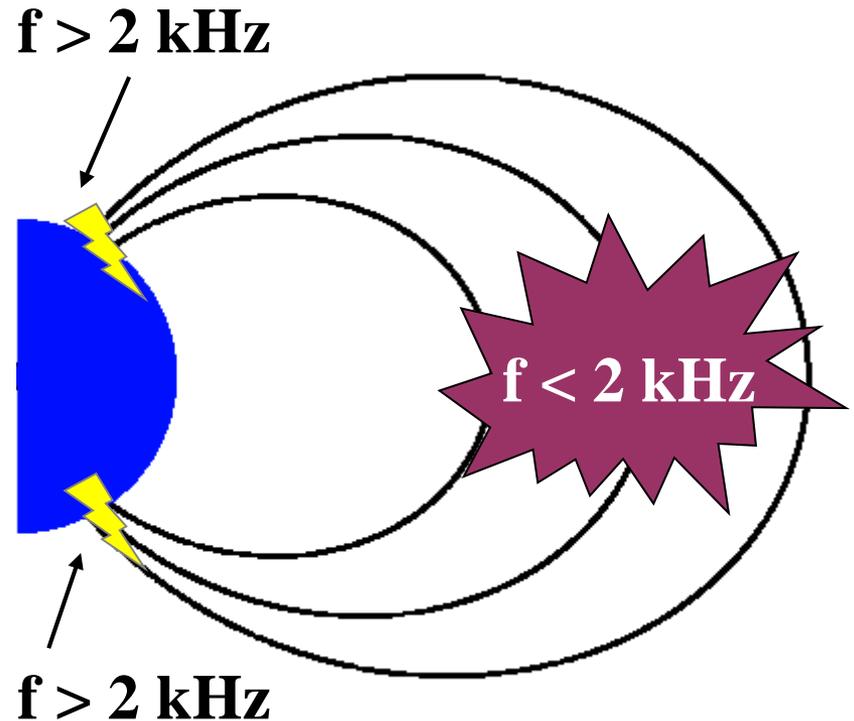
L	τ (days)
3.0	8.0 ± 2.6
2.5	3.6 ± 1.6
2.0	20 ± 9.5

Lifetime of 2 MeV Electrons



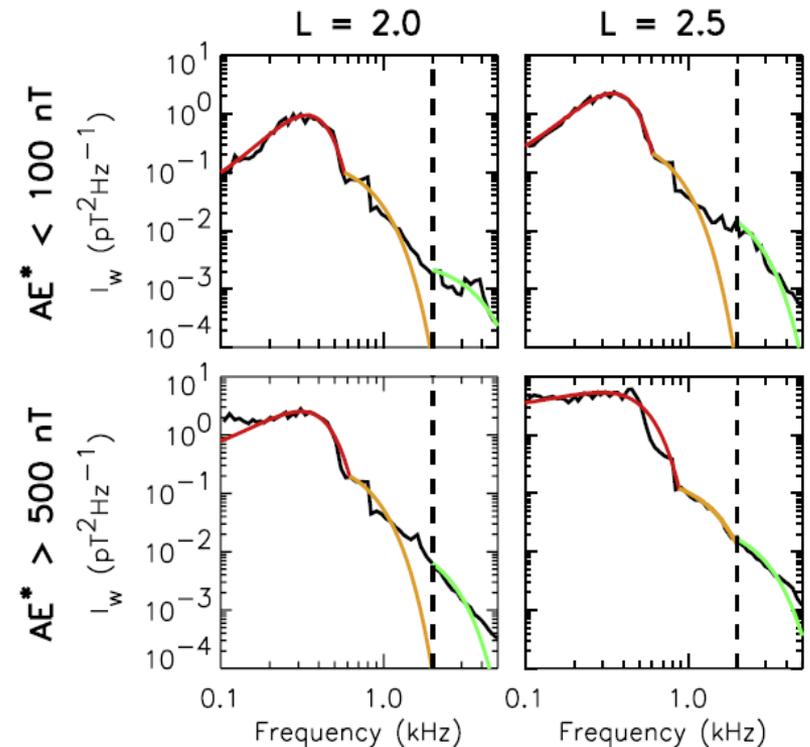
Broadband Plasmaspheric Emissions

- Broadband plasmaspheric emissions can be split into two categories [Meredith *et al.*, 2006]:
 - Plasmaspheric hiss
 - $100 \text{ Hz} < f < 2 \text{ kHz}$
 - generated by whistler mode chorus
 - Lightning-generated whistlers
 - $2 \text{ kHz} < f < 5 \text{ kHz}$
 - produced by thunderstorms on Earth



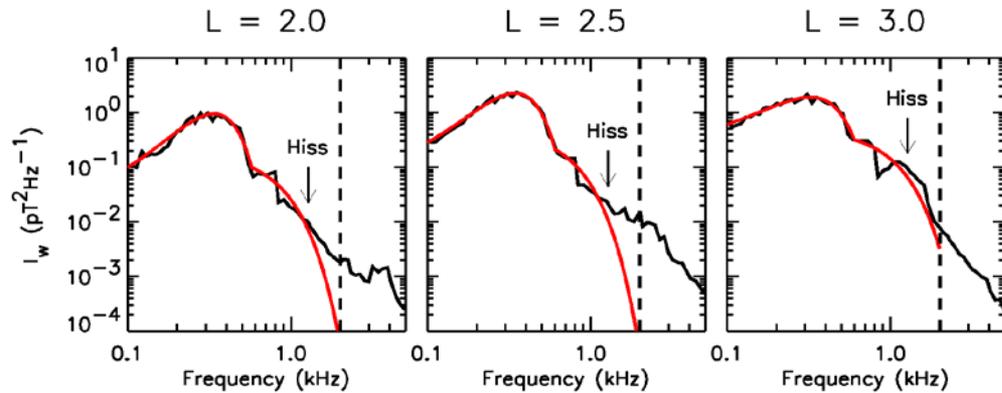
Calculation of Loss Timescales from Waves

- Use global models of the wave spectral intensity based on CRRES observations.
- Calculate bounce-averaged pitch angle diffusion rates using the PADIE code [Glauert and Horne, 2005].
- Determine the loss timescale and the evolution of the pitch angle distribution using the 1D pitch angle diffusion equation following Lyons et al., [1972].



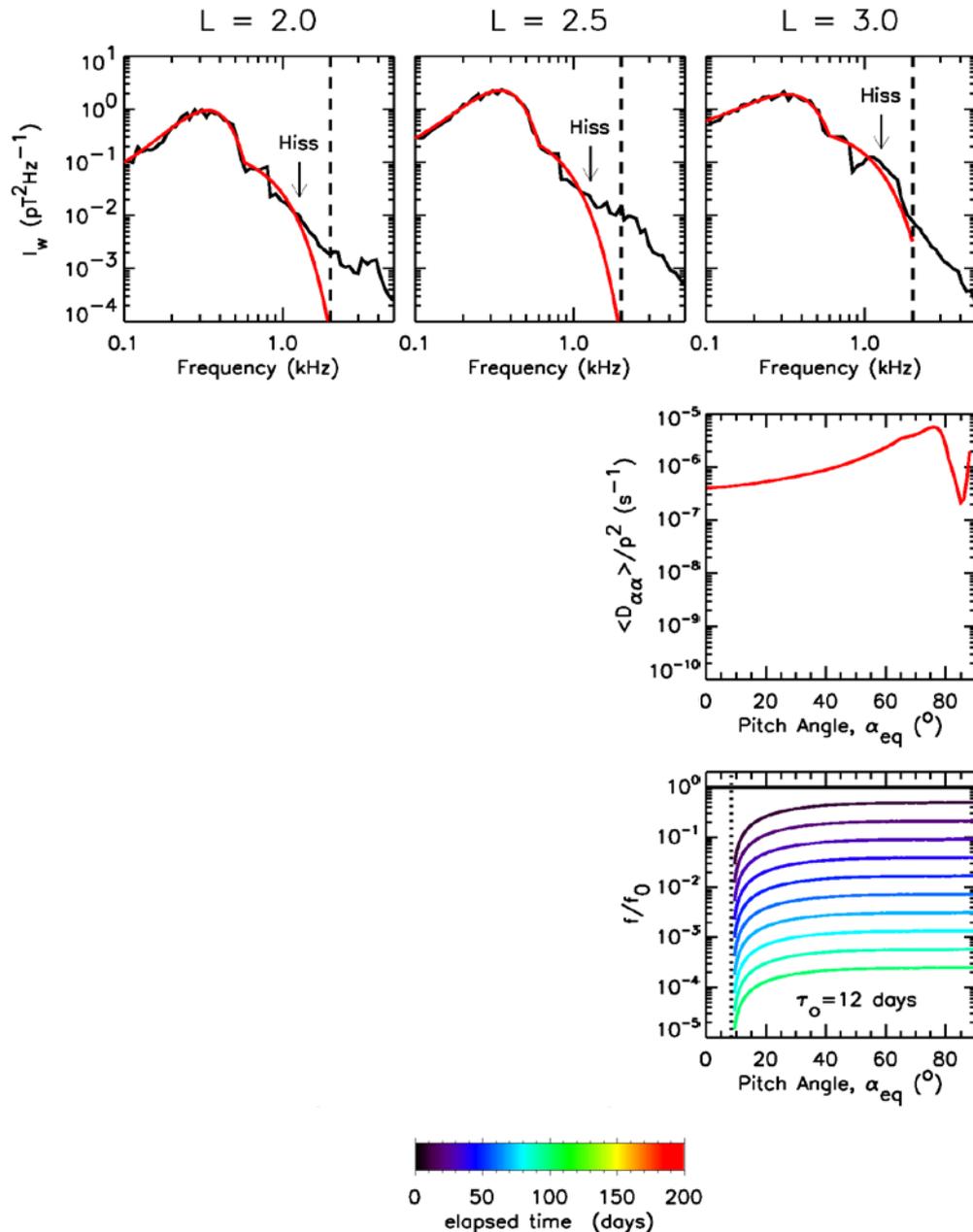
Meredith et al., JGR, 2007

Lifetimes due to Hiss



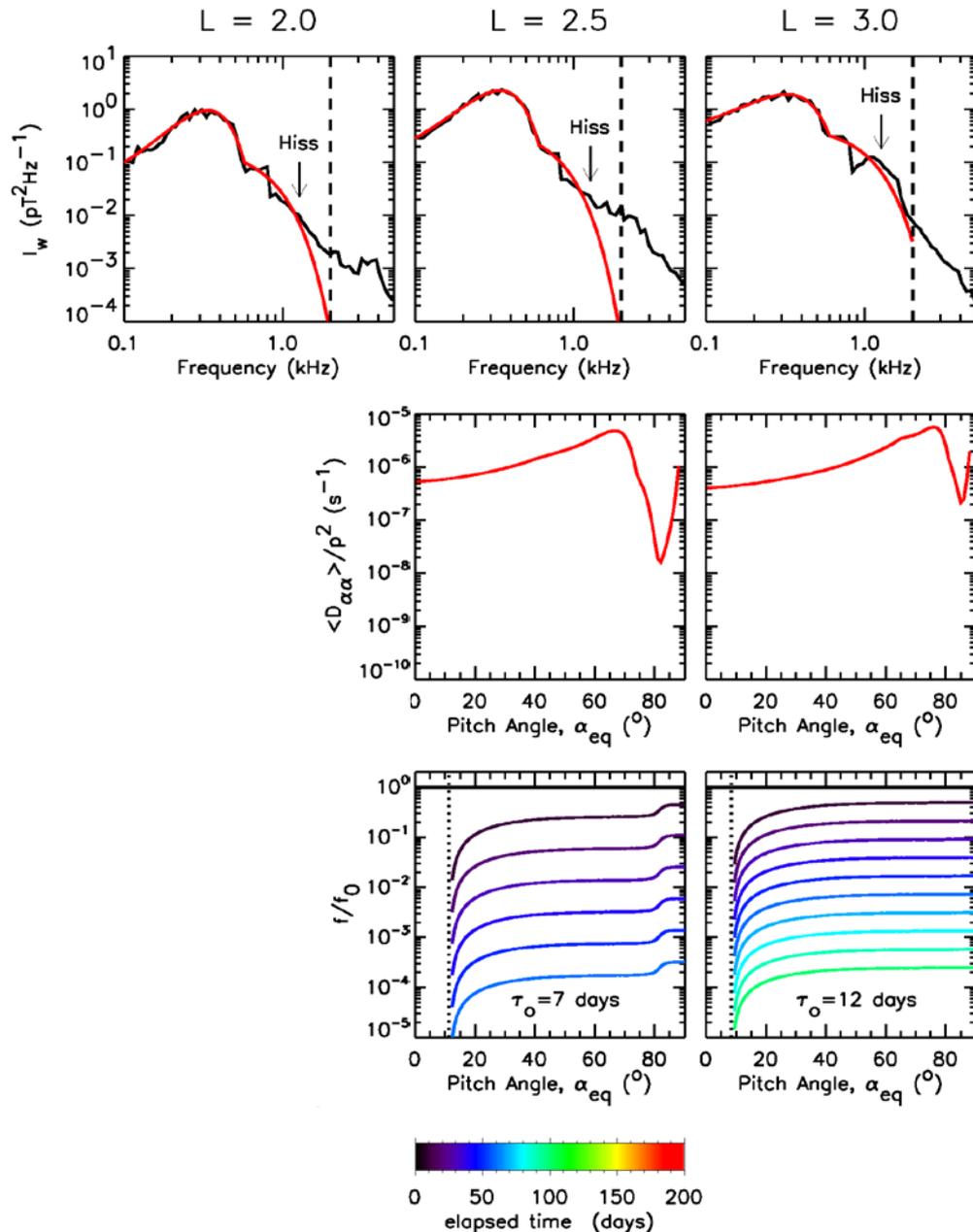
- Wave power in the frequency range 100 Hz – 5 kHz is dominated by plasmaspheric hiss.
- We begin by considering the effect of hiss during quiet conditions ($AE^* < 100$ nT).

Lifetimes due to Hiss



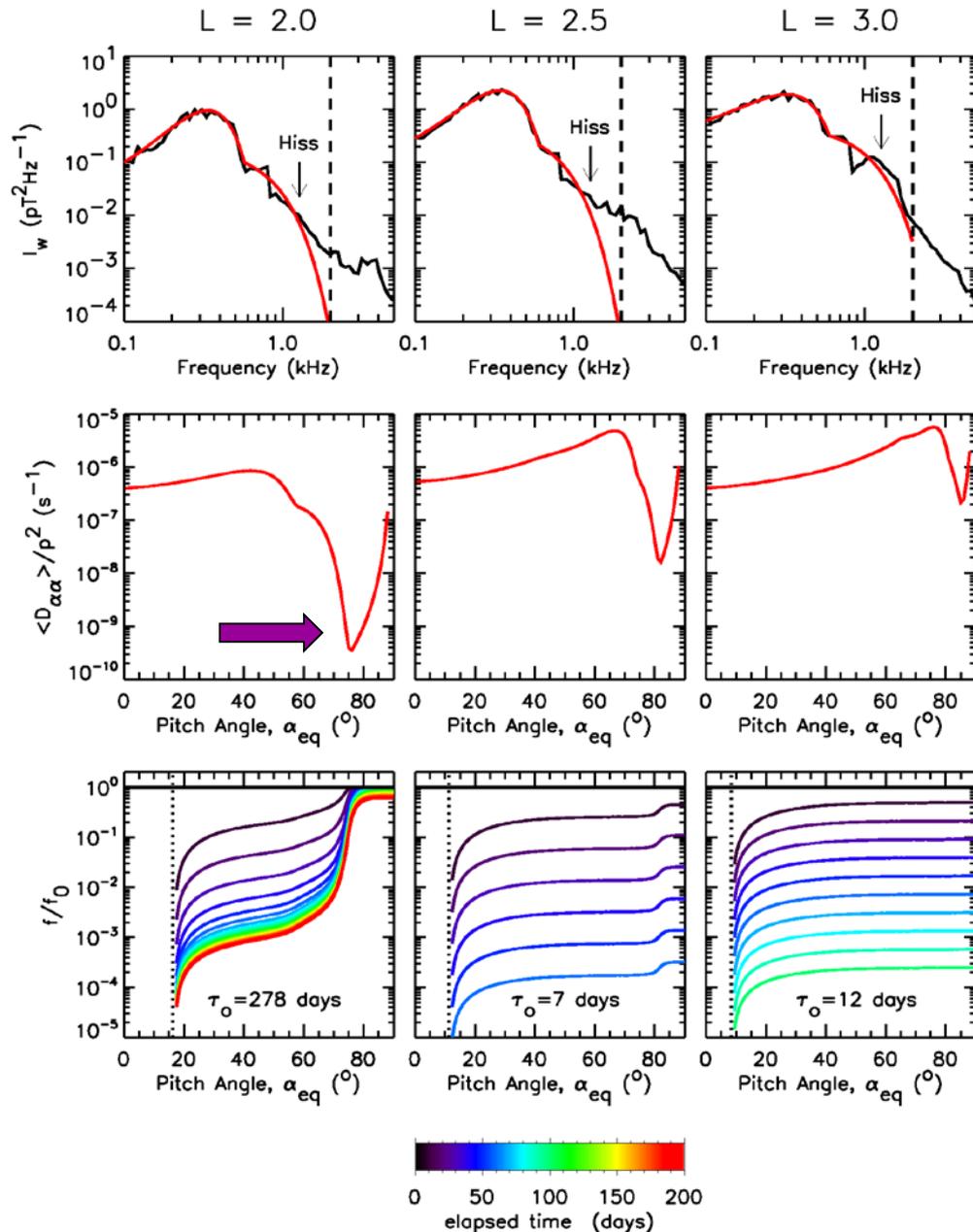
- At $L = 3.0$ there is a minimum in the diffusion rates at large pitch angles, but the diffusion rate does not fall by more than a factor of 2 compared to the edge of the loss cone.
- The pitch angle distribution quickly reaches an equilibrium state and decays exponentially at all pitch angles on a timescale of 12 days.

Lifetimes due to Hiss



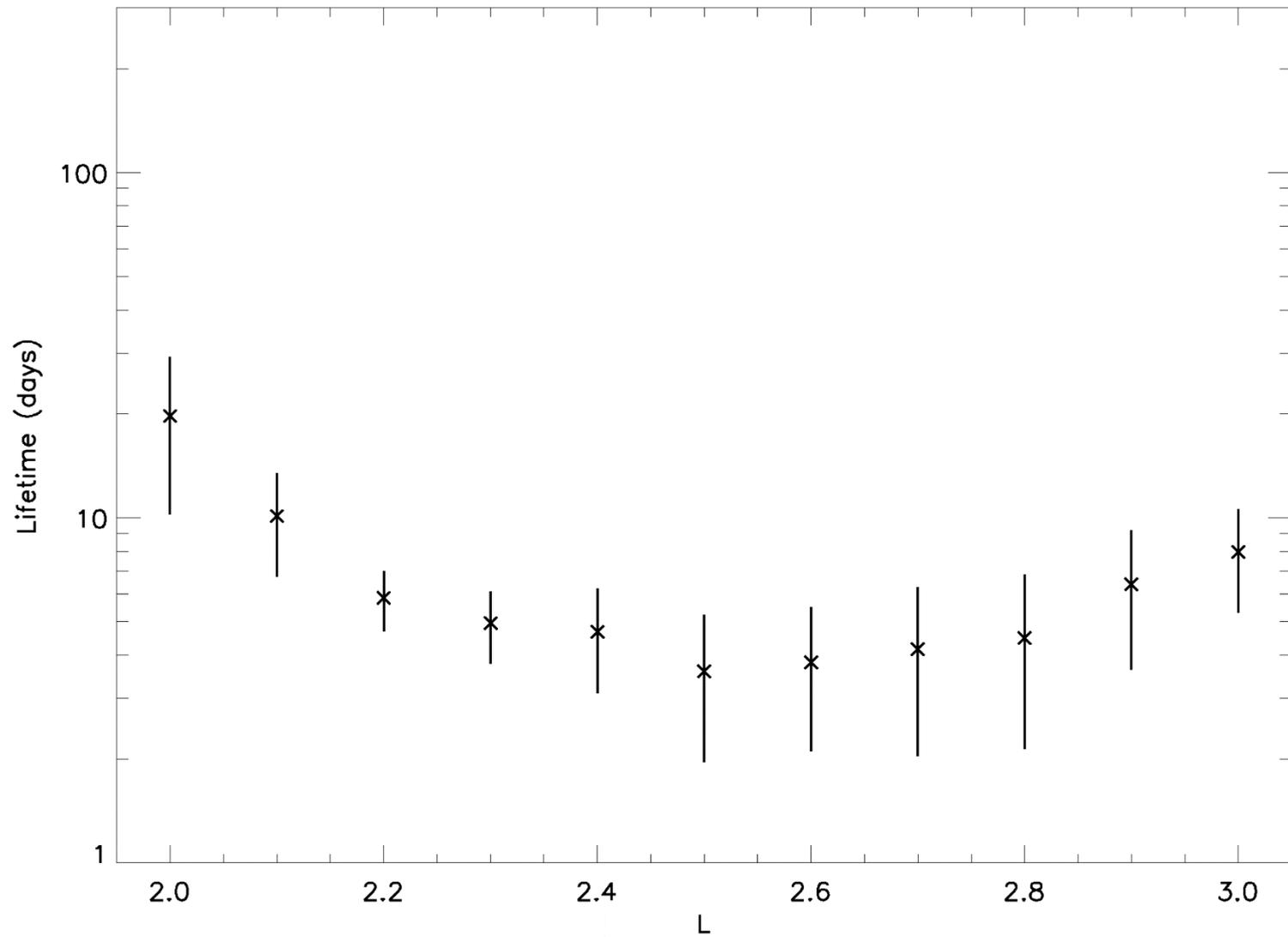
- At $L = 2.5$ there is a larger minimum in the diffusion rate.
- The pitch angle distribution quickly reaches an equilibrium state and decays exponentially at all pitch angles on a timescale of 7 days.

Lifetimes due to Hiss

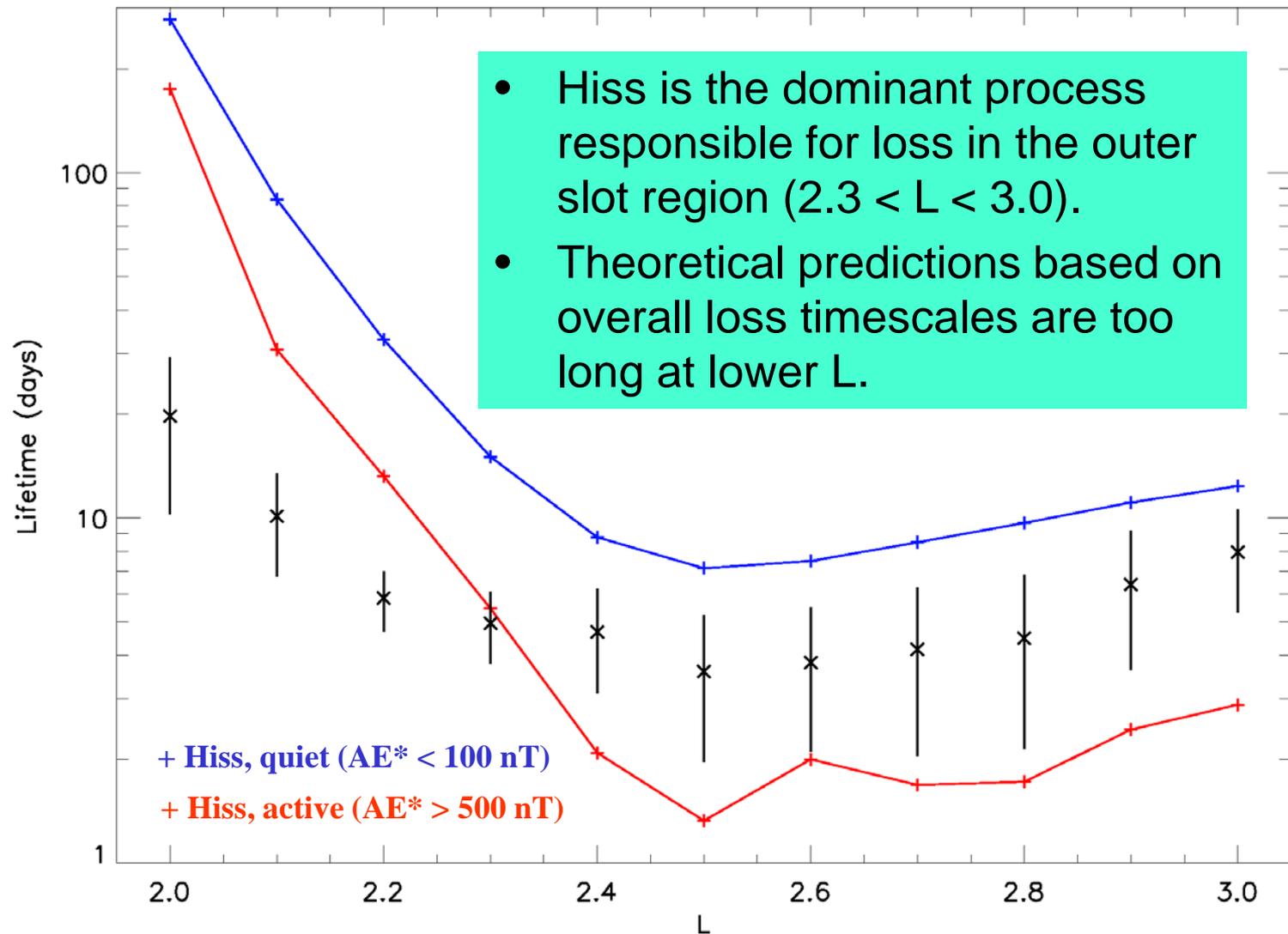


- At $L = 2.0$ there is a very deep minimum in the diffusion rate.
- This dramatically effects the evolution of the PAD:
 - The decay is pitch angle dependent.
 - The distribution initially decays more rapidly at smaller pitch angles.
- Once an equilibrium shape is reached the entire distribution decays with a timescale of 278 days.

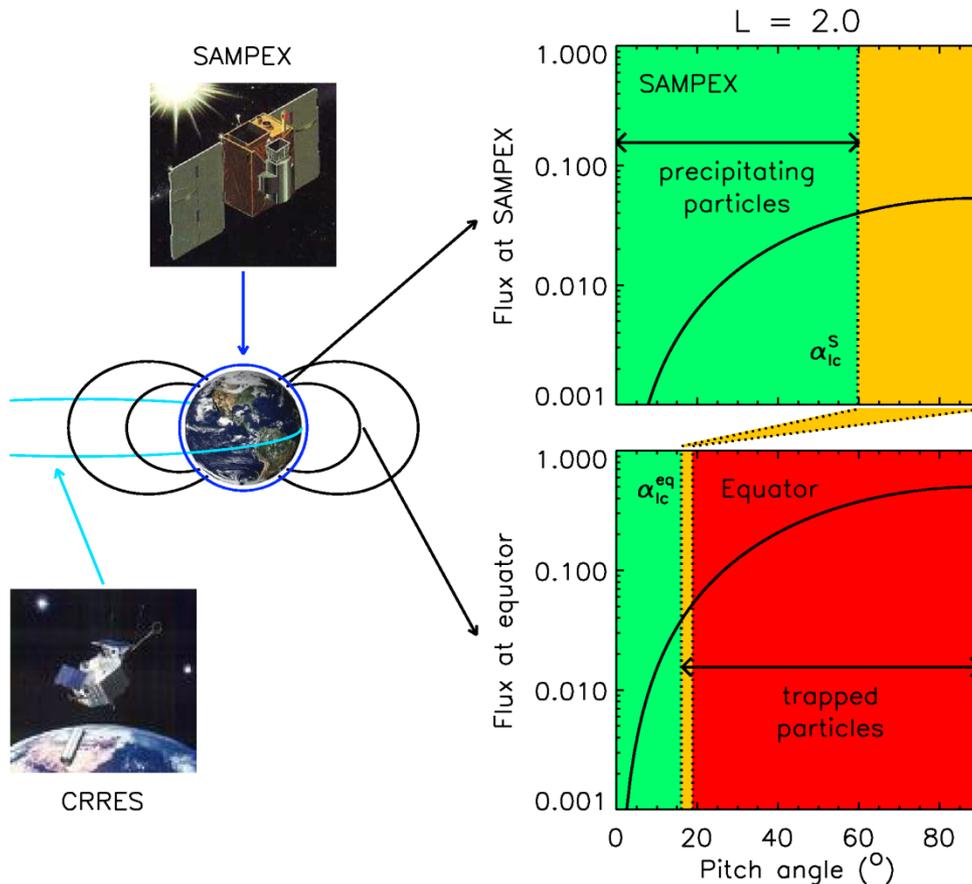
Lifetime of 2-6 MeV Electrons



Lifetimes due to Plasmaspheric Hiss

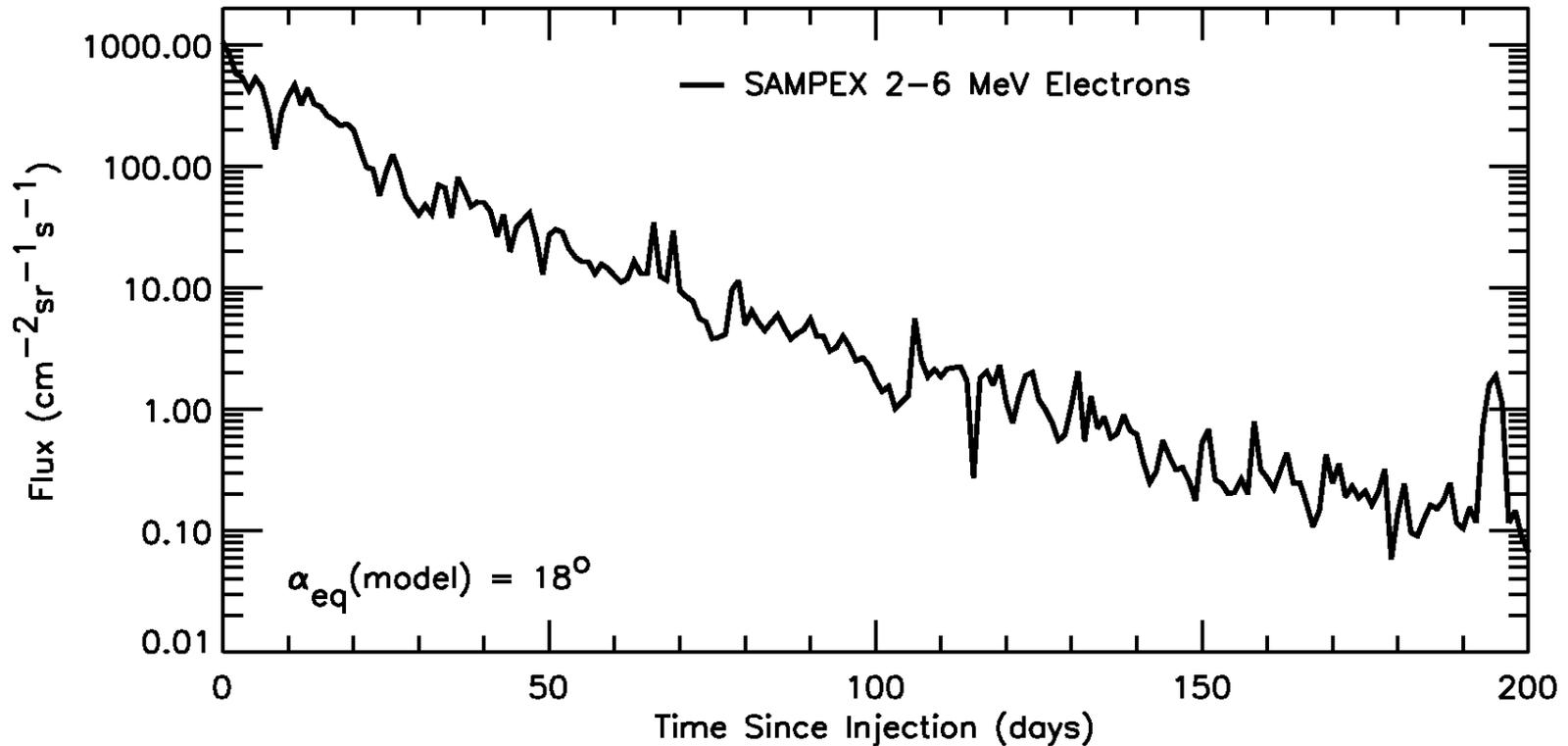


SAMPEX Measurements at L = 2.0



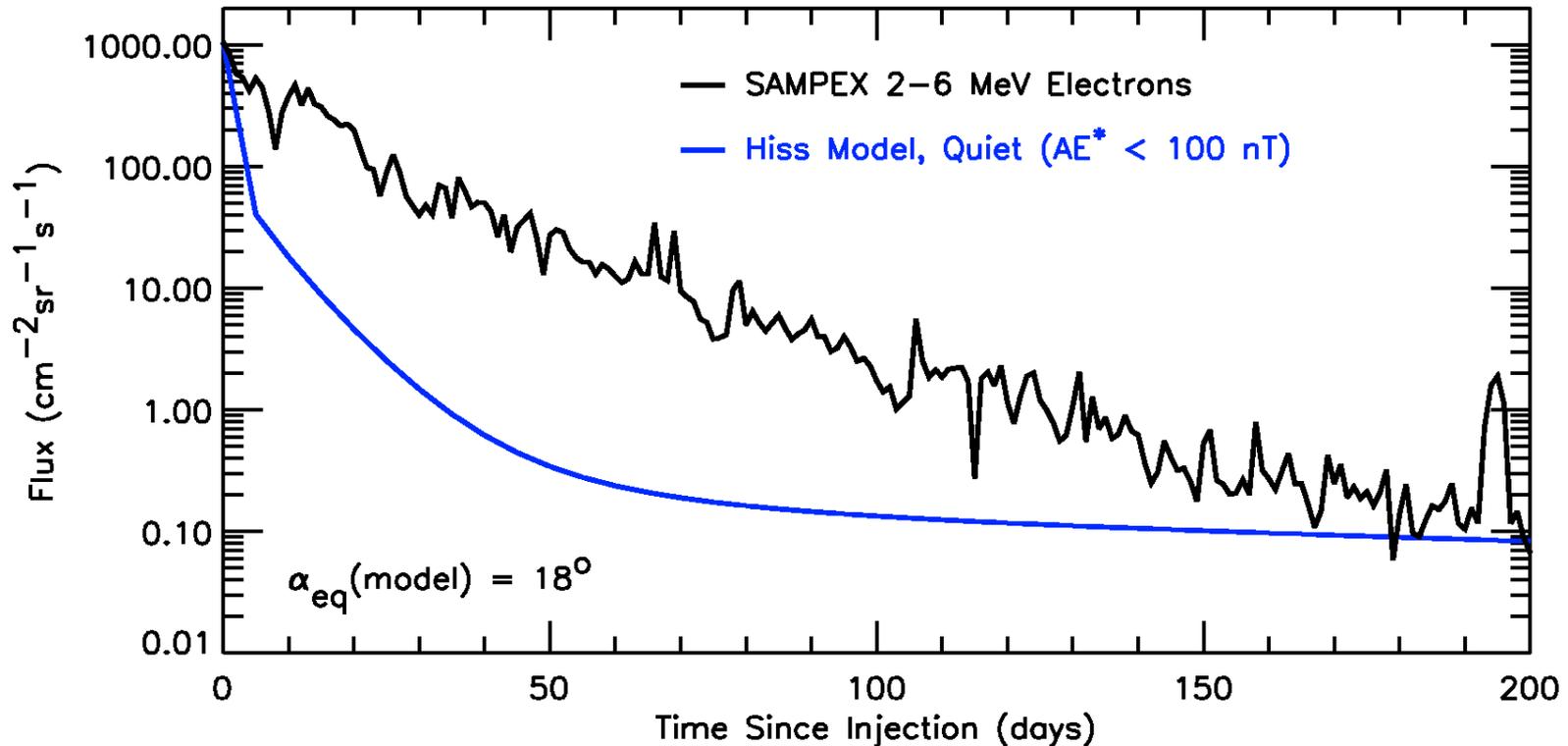
- At L = 2.0 the decay is pitch angle dependent.
- SAMPEX makes measurements at low altitudes (~600 km).
- A locally mirroring particle has an equatorial pitch angle of $\sim 18^\circ$ at L = 2.0.
- Here we should compare SAMPEX measurements with the model flux at 18° .

Flux Decay at L = 2.0

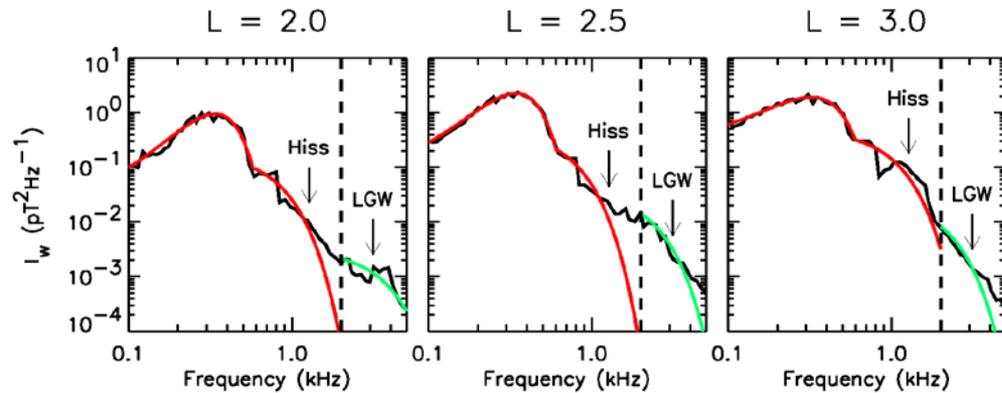


- SAMPEX measures a relatively constant exponential decay.

Flux Decay at L = 2.0



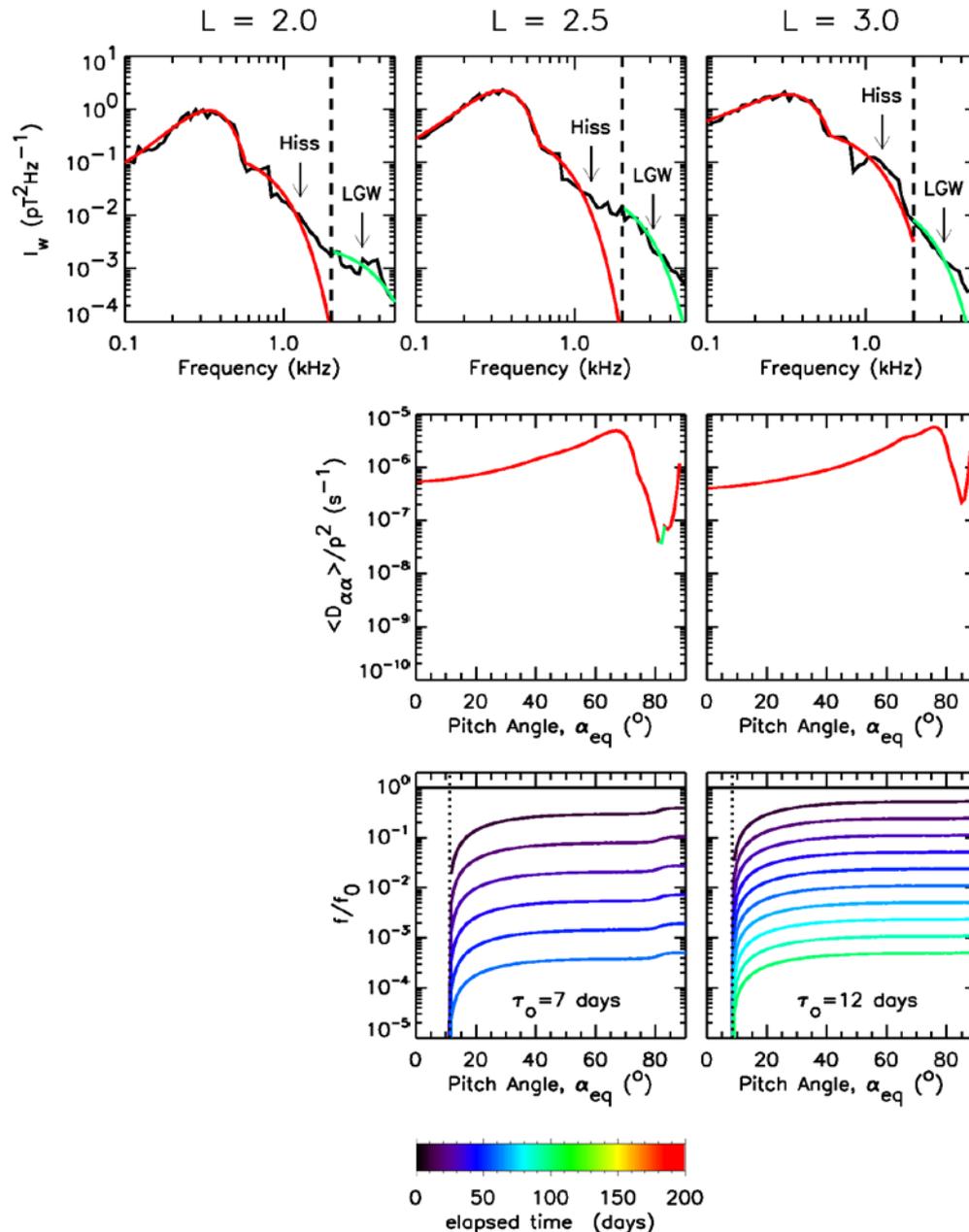
- Decay of the model flux at $\alpha_{\text{eq}} = 18^\circ$ is time-dependent and cannot explain the observations.
- Losses due to plasmaspheric hiss alone cannot explain the losses at L = 2.0.



Losses due to Hiss and LGWs

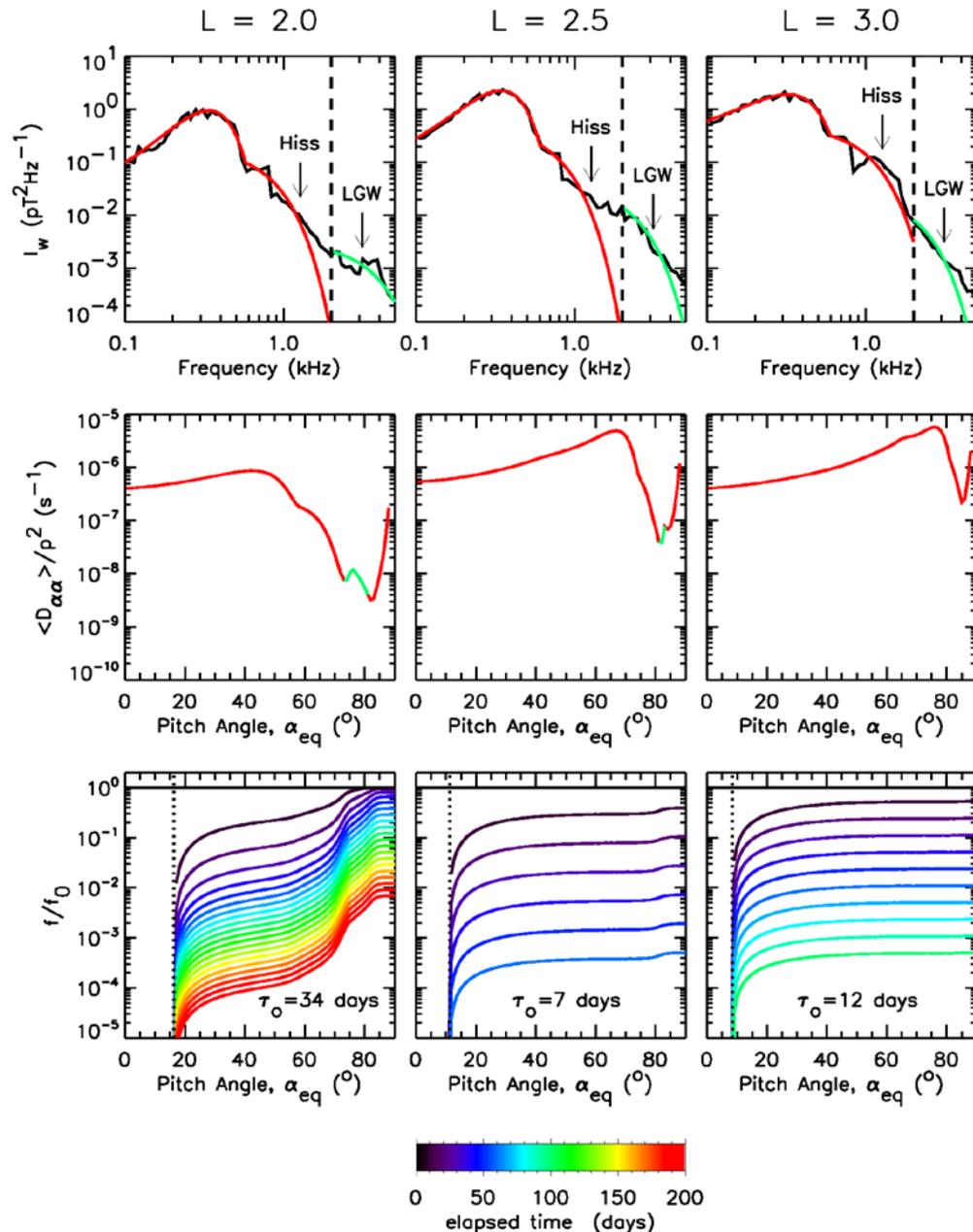
- We now consider the combined effect of lightning-generated whistlers (LGWs) and plasmaspheric hiss.
- LGWs dominate the spectrum above 2 kHz and so we add an additional component from 2-5 kHz.
- We again present the results for quiet conditions ($AE^* < 100$ nT)

Losses due to Hiss and LGWs



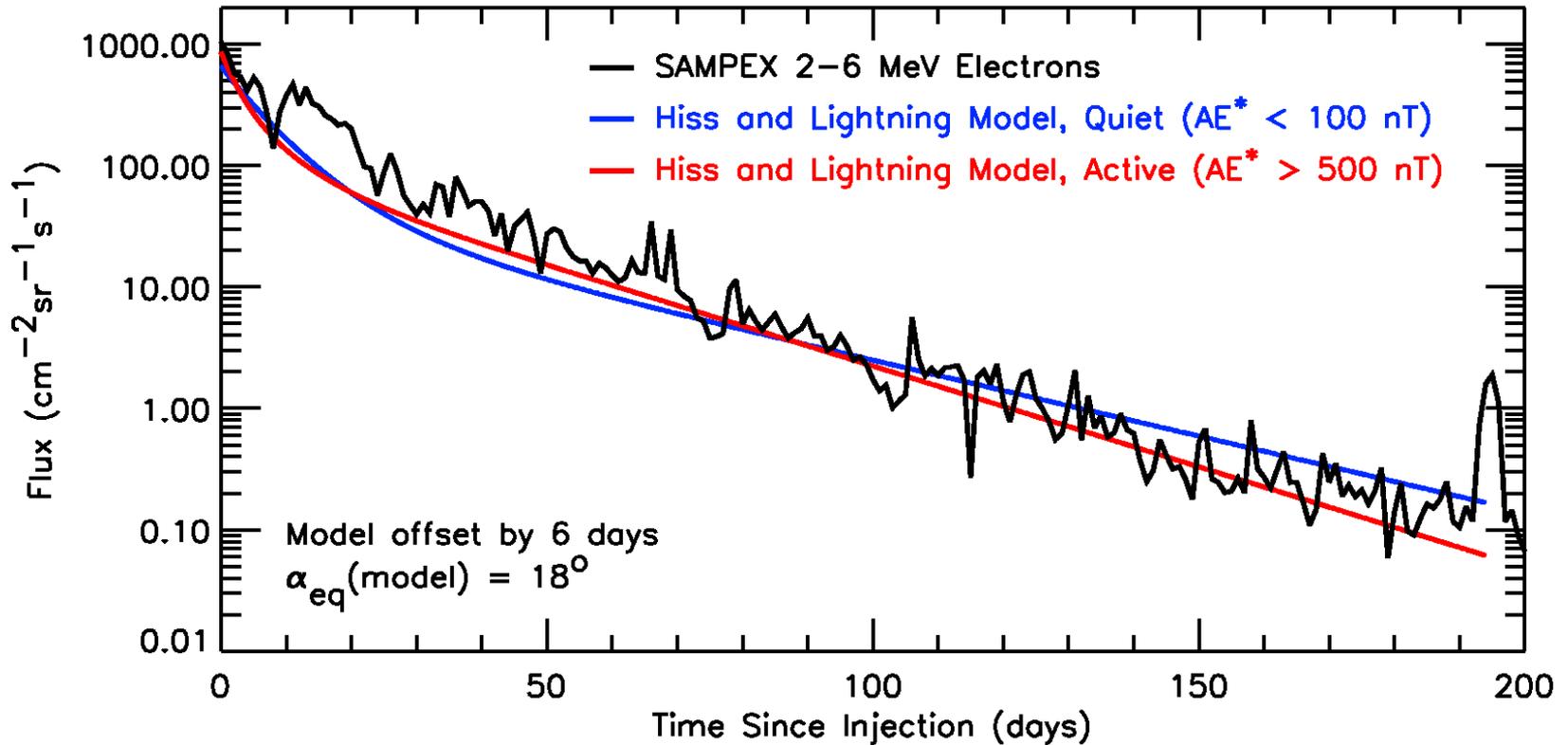
- At $L = 2.5$ and $L = 3.0$ there is little or no change in the diffusion rates and virtually no change in the loss rates.

Losses due to Hiss and LGWs



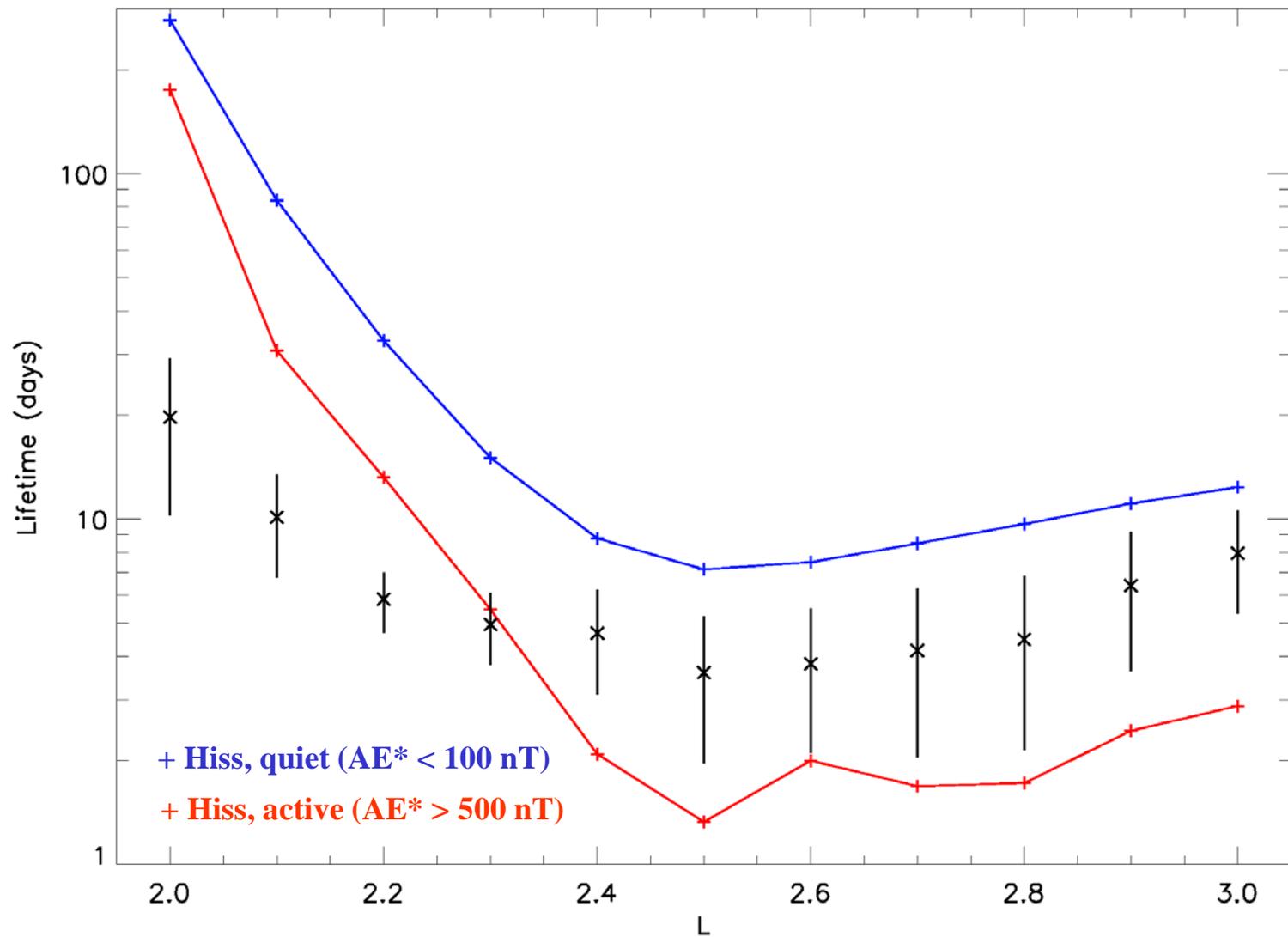
- At $L = 2.0$, the effect of the additional wave power is to increase the diffusion rates in the deep minimum.
- The distribution now evolves more quickly to an equilibrium state and decays with a lifetime of 34 days.

Evolution of the Flux at SAMPEX altitudes

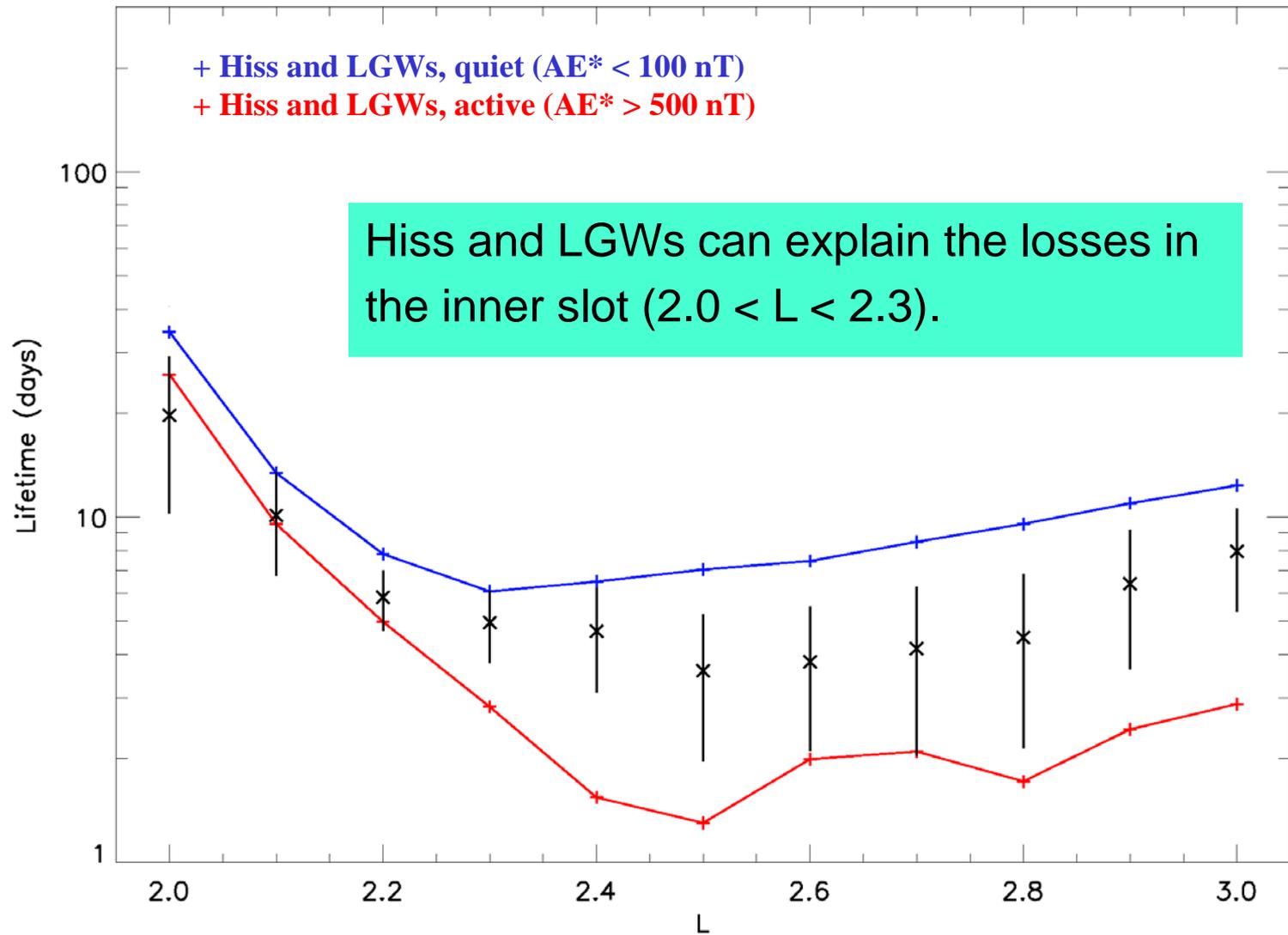


- Decay of the model flux at $\alpha_{\text{eq}} = 18^\circ$ gives reasonable agreement with the data.

Lifetimes due to Plasmaspheric Hiss

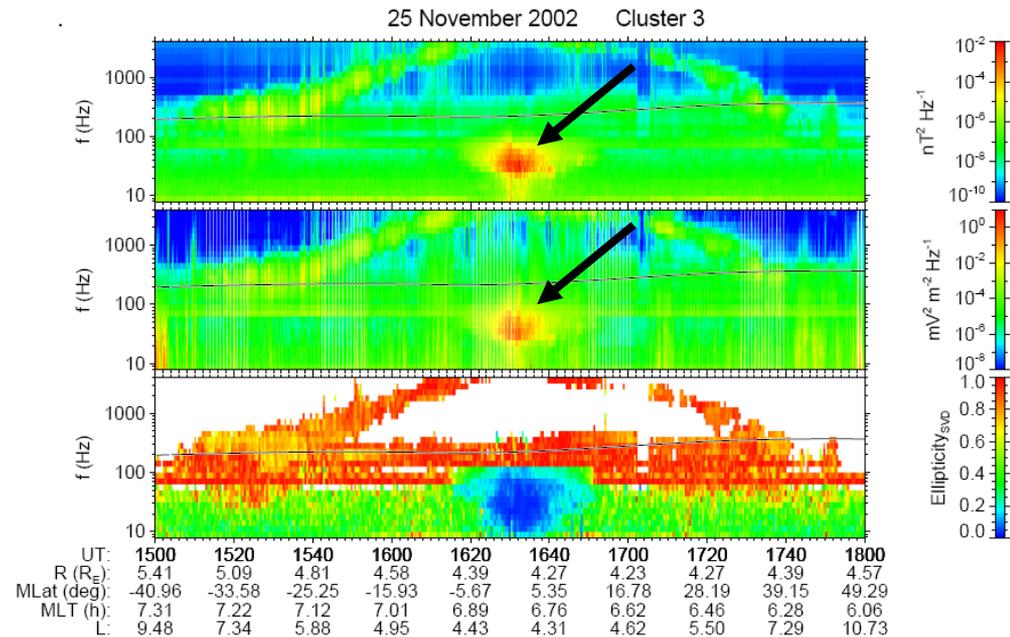


Lifetimes due to Hiss and LGWs



Magnetosonic Waves

- Magnetosonic waves may also scatter radiation belt electrons at large pitch angles [Horne *et al.*, 2007].
- These waves are
 - intense electromagnetic emissions
 - $f_{\text{CH}} < f < f_{\text{LHR}}$
 - observed near the equatorial plane
 - $2 < L < 7$

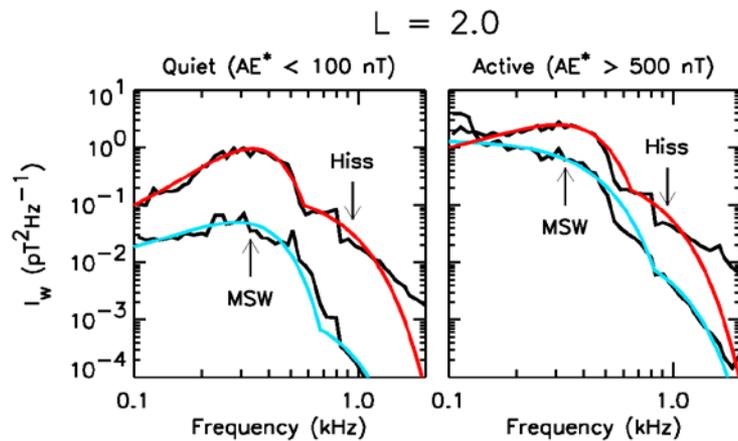


Horne *et al.*, 2007



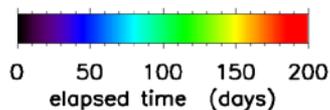
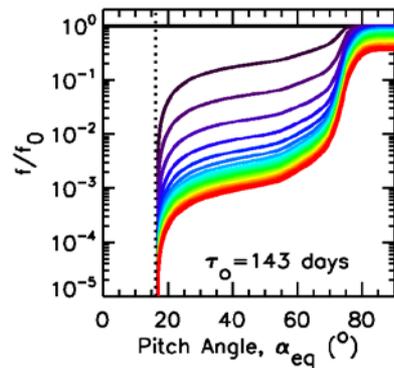
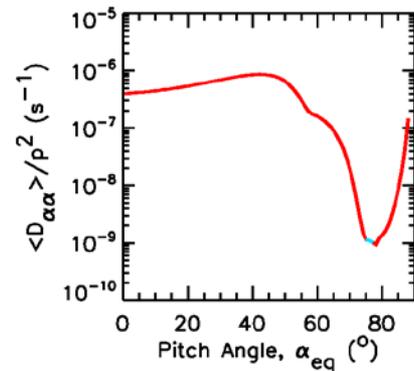
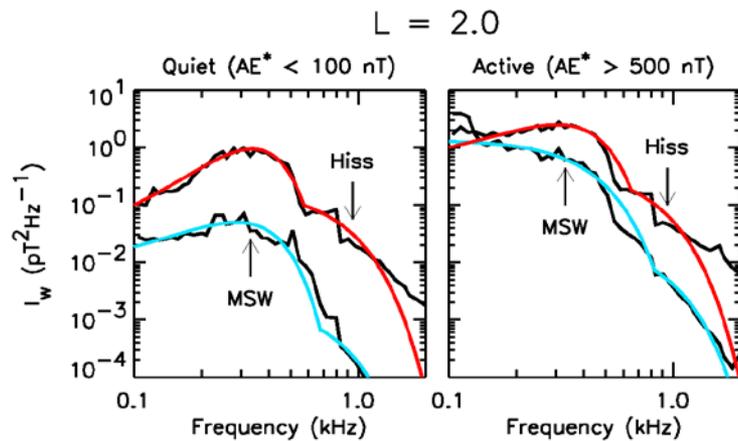
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Losses due to MSWs

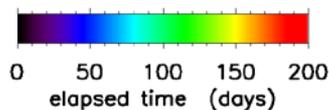
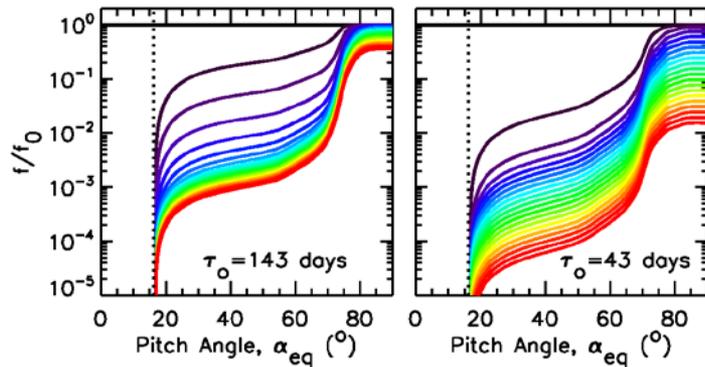
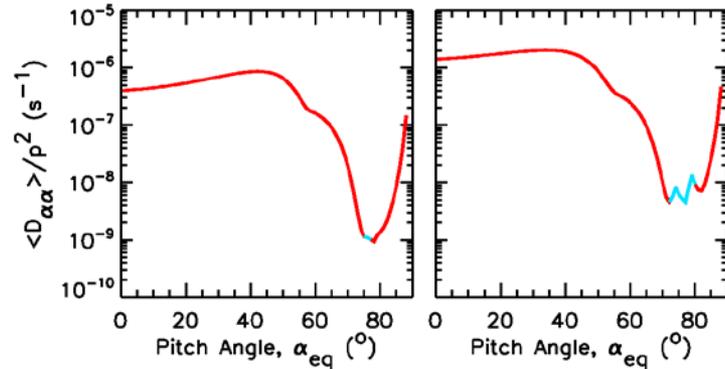
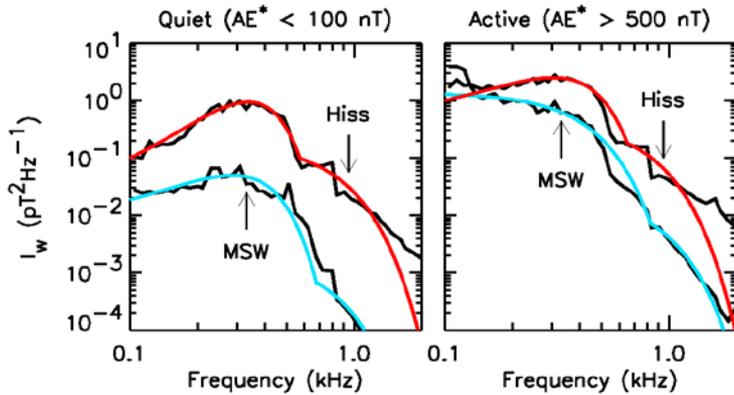
- We cannot uniquely identify MSWs in the CRRES data.
- We use global averages of the wave spectral intensity observed within $\pm 3^\circ$ of the equator as an estimate of the upper limit.



Losses due to MSWs

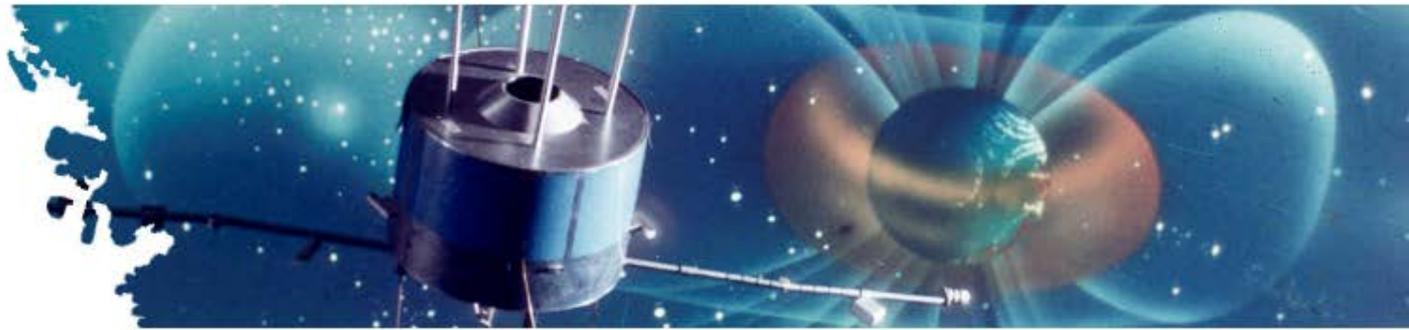
- During quiet conditions MSWs make only a small change to the diffusion rates at large pitch angles.
- The resulting lifetime of 143 days is too long to explain the observations.

$L = 2.0$



Losses due to MSWs

- During active conditions MSWs increase the diffusion rates at large pitch angles.
- The resulting loss timescale is 43 days.
- MSWs could contribute to the loss timescales during active conditions.
- Better models of MSWs required to establish the role of these waves.



Conclusions

- Pitch angle scattering by plasmaspheric hiss is responsible for electron loss in the outer slot region ($2.3 < L < 3.0$)
- In the inner slot region ($2.0 < L < 2.3$) electron loss is driven by plasmaspheric hiss and lightning-generated whistlers.
- Magnetosonic waves may also contribute to electron loss at $L = 2.0$ but better wave models are required to determine the precise role of these waves.



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Loss Timescales for $2.0 < L < 4.0$

Electron Loss Timescale Versus Energy

